

FRACTIONAL SUPERSYMMETRIC- $sl(2)$

Yasemen UCAN*

*Yıldız Teknik Üniversitesi, Kimya-Metalurji Fakültesi, Matematik Mühendisliği Bölümü,
Davutpaşa-İSTANBUL*

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KESİRSEL SÜPERSİMETRİK- $sl(2)$

ÖZET

Permutasyon grubunun S_3 invariant formları üzerinde kurulan Lie cebirinin kübik kökü Hopf cebri formalizminde ifade edildi. $n=3$ 'te $sl(2)$ 'nin $N=4$ kesirsel süper genellemesini gözönüne aldık.

Anahtar Sözcükler: Kesirsel süpercebirlere, $sl(2)$ Lie cebiri, Kesirsel süper- $sl(2)$

ABSTRACT

The 3rd root of Lie algebra based on the permutation group S_3 invariant forms is formulated in the Hopf algebra formalism. We consider $N=4$ fractional super generalizations of $sl(2)$ at $n=3$

Keywords: Fractional superalgebras, $sl(2)$ Lie algebra, Fractional super- $sl(2)$

1. INTRODUCTION

To arrive at a superalgebra one adds new elements Q_a to generators X_j of the corresponding Lie algebra and defines the relations

$$\{Q_a, Q_b\} = b_{ab}^j X_j \quad (1)$$

observing that the anticommutator in the above relation is invariant under the cyclic Z_2 or permutation S_2 groups anticommutator. To arrive at cubic root of a Lie algebra g , instead of (1) has the cubic relation

* e-mail: ucan@yildiz.edu.tr ; tel: (0212) 449 1762

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$$Q_a Q_b Q_g + Q_g Q_a Q_b + Q_b Q_g Q_a = b_{abg}^j X_j \quad (2)$$

which is Z_3 invariant and the cubic relation

$$Q_a \{Q_b, Q_g\} + Q_b \{Q_a, Q_g\} + Q_g \{Q_a, Q_b\} = b_{abg}^j X_j \quad (3)$$

which is S_3 invariant. From the above relations only (3) appears to be consistent at the co-algebra level. So we used the relation (3).

Fractional superalgebras based on S_n invariant form were first introduced in [1,2] and later constructed in the Hopf algebra context and defined their dual in [3]. In this paper, according to [3] we discuss fractional super- $sl(2)$ for $N=4$.

There are other approaches to fractional supersymmetry in the Literature [4-9]. For example, one can arrive at fractional super algebras by using quantum groups at the roots of unity [10]. The plan of the paper is as follows. In the section 2, we give a formulation of fractional superalgebras in the Hopf algebra formalism from the [3]. In the section 3, we consider $N=4$ fractional supergeneralization of $sl(2)$ at $n=3$. we denoted this algebra by $U_3^4(sl(2))$.

2. REVIEW OF FRACTIONAL SUPERALGEBRAS

Let $U(g)$ be the universal enveloping algebra of a Lie algebra g generated by X_j $j=1,2,\dots, \dim(g)$ with

$$[X_i, X_j] = \sum_{k=1}^{\dim(g)} c_{ij}^k X_k \quad (4)$$

Where c_{ij}^k are the structure constants of the Lie algebra g . The Hopf algebra structure of $U(g)$ is given by

$$\Delta(X_j) = X_j \otimes 1 + 1 \otimes X_j, \quad e(X_j) = 0, \quad S(X_j) = -X_j. \quad (5)$$

To arrive at cubic root of $U(g)$. we shall use S_3 invariant form. Therefore, we defined an algebra generated by X_j , $j=1,\dots, \dim(g)$ and Q_a, K , $a=1,\dots, N$ satisfying the relations (4) and

$$\{Q_a, Q_b, Q_g\} = b_{abg}^j X_j \quad (6)$$

$$[Q_a, X_j] = a_{ab}^j Q_b \quad (7)$$

and

$$K Q_a = q Q_a K, \quad q^3 = 1, \quad K^3 = 1 \quad (8)$$

where

$$\{Q_a, Q_b, Q_g\} \equiv Q_a \{Q_b, Q_g\} + Q_b \{Q_a, Q_g\} + Q_g \{Q_a, Q_b\}$$

is the S_3 invariant form, c_{ij}^k and a_{ab}^j, b_{abg}^j are the structure coefficients satisfying the Jacobi and super Jacobi identities. This algebra is denoted by the symbol $U_3^N(g)$. The above algebra is a Hopf algebra with the following co structures [3]:

$$\Delta(Q_a) = Q_a \otimes 1 + K \otimes Q_a, \quad \Delta(K) = K \otimes K, \tag{9}$$

$$e(Q_j) = 0, \quad e(K) = 1, \tag{10}$$

$$S(Q_j) = -K^2 Q_j, \quad S(K) = K^2. \tag{11}$$

To define structure constant a_{ab}^j and b_{abg}^j we have to derive identities involving the commutator and S_3 invariant form. One can check that relations

$$[A, [B, C]] + [C, [A, B]] + [B, [C, A]] = 0 \tag{12}$$

$$[A, \{B, C, D\}] + \{[B, A], C, D\} + \{B, [C, A], D\} + \{B, C, [D, A]\} = 0 \tag{13}$$

and

$$[A, \{B, C, D\}] + [B, \{A, C, D\}] + [C, \{B, A, D\}] + [D, \{B, C, A\}] = 0 \tag{14}$$

are identically satisfied [3]. The relation (12) is the usual Jacobi identity.

Inserting

$$A = X_i, \quad B = X_j, \quad C = Q_a \tag{15}$$

into (12) and using (7) and (4) we get

$$\sum_{s=1}^N (a_{as}^i a_{sb}^j - a_{as}^j a_{sb}^i) = \sum_{k=1}^{\dim(g)} c_{ij}^k a_{ab}^k \tag{16}$$

Comparing the above relation with (4) we conclude that the $N \times N$ matrices $a^j \equiv (a_{ab}^j)_{a,b=1}^N$ define a N-dimensional representation of a given Lie algebra. Of course, these matrices are not unique.

Let us now consider restrictions on structure coefficients coming from the other identities. Inserting

$$A = X_k, \quad B = Q_a, \quad C = Q_b, \quad D = Q_g \tag{17}$$

into the identity (13) we get

$$\sum_{s=1}^N (a_{as}^k b_{sbg}^i + a_{bs}^k b_{sag}^i + a_{gs}^k b_{sba}^i) = \sum_{j=1}^{\dim(g)} c_{jk}^i b_{abg}^j \tag{18}$$

and

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$$A = Q_s \quad , \quad B = Q_a \quad , \quad C = Q_b \quad , \quad D = Q_g \quad (19)$$

into (14) and using (6), (7) we obtained following the relation

$$\sum_{k=1}^{\dim(g)} (b_{abg}^k a_{st}^k + b_{sab}^k a_{gt}^k + b_{gsa}^k a_{bt}^k + b_{bgs}^k a_{at}^k) = 0 \quad (20)$$

3. N=4 FRACTIONAL SUPER- $sl(2)$

We know that the generators of the algebra $sl(2)$ satisfy the following commutation relations

$$[X_1, X_2] = X_3 \quad [X_3, X_1] = 2X_1 \quad [X_3, X_2] = -2X_2 \quad (21)$$

From the relation (21) one has

$$c_{12}^3 = 1 \quad c_{31}^1 = 2 \quad c_{32}^2 = -2 \quad (22)$$

For N=4, the matrix $a^j = \{a_{ab}^j\}$ due to (16) is an arbitrary 4-dimensional representation of $sl(2)$. The solution of (18) and (20) for b_{abg}^j is fully determined by this representation. where b_{abg}^j is symmetric in a, b and g . We consider N=4 super

generalization of $sl(2)$ at n=3, that is $q = e^{i\frac{p}{3}}$. We have different superalgebras depending on the choice of a^j .

(i) we take $a_{ab}^j = 0$. Then the relations (18) and (20) imply $b_{abg}^j = 0$. The obtained structure constants imply that the fractional superalgebra $U_3^4(sl(2))$ is the direct product of the universal enveloping algebra $U(sl(2))$ with the Hopf algebra generated by Q_1, Q_2, Q_3, Q_4 and K satisfying the relations

$$K Q_a = q Q_a K \quad \{Q_a, Q_b, Q_g\} = 0 \quad K^3 = 1 \quad (23)$$

and the Hopf algebra structure (9)- (11).

(ii) Take the vector representation

$$a^1 = \begin{pmatrix} 0 & \sqrt{3} & 0 & 0 \\ 0 & 0 & 2 & 0 \\ 0 & 0 & 0 & \sqrt{3} \\ 0 & 0 & 0 & 0 \end{pmatrix} \quad a^2 = \begin{pmatrix} 0 & 0 & 0 & 0 \\ \sqrt{3} & 0 & 0 & 0 \\ 0 & 2 & 0 & 0 \\ 0 & 0 & \sqrt{3} & 0 \end{pmatrix} \quad a^3 = \begin{pmatrix} 3 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -3 \end{pmatrix} \quad (24)$$

The substitution of

$$a_{12}^1 = a_{21}^2 = a_{34}^1 = a_{43}^2 = \sqrt{3} \quad a_{23}^1 = a_{32}^2 = 2 \quad (25)$$

$$a_{11}^3 = 3, \quad a_{22}^3 = 1, \quad a_{33}^3 = -1, \quad a_{44}^3 = -3$$

into (20) and (18), will give all parameters b_{abg}^j are zero. Thus we obtained the following fractional superalgebra :

$$\{Q_a, Q_b, Q_g\} = 0 \tag{26}$$

and

$$[Q_1, X_1] = \sqrt{3} Q_2 \quad [Q_1, X_3] = 3Q_1$$

$$[Q_2, X_1] = 2Q_3 \quad [Q_2, X_2] = \sqrt{3} Q_1 \quad [Q_2, X_3] = Q_2 \tag{27}$$

$$[Q_3, X_1] = \sqrt{3} Q_4 \quad [Q_3, X_2] = 2Q_2 \quad [Q_3, X_3] = -Q_3$$

$$[Q_4, X_2] = \sqrt{3} Q_3 \quad [Q_4, X_3] = -3Q_1$$

Note that, for N=3 the relations (26) are not zero [3].

(iii) Assume that two of the fractional super generators Q_1, Q_2, Q_3 and Q_4 transform as spinors and the remaining two transforms as scalars, that is

$$a^1 = \begin{bmatrix} 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix} \quad a^2 = \begin{bmatrix} 0 & 0 & 0 & 0 \\ 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix} \quad a^3 = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix} \tag{28}$$

The substitution of

$$a_{12}^1 = a_{21}^2 = a_{11}^3 = 1 \quad a_{22}^3 = -1 \tag{29}$$

into (20) gives

$$\begin{aligned} b_{222}^1 &= -3b_{112}^2 = 3b_{122}^3 \\ b_{122}^1 &= -\frac{1}{3}b_{111}^2 = b_{112}^3 \\ b_{223}^1 &= -b_{113}^2 = 2b_{123}^3 \\ b_{224}^1 &= -b_{114}^2 = 2b_{124}^3 \end{aligned} \tag{30}$$

The substituting these into (18) one finds that the only solution is $b_{abg}^j = 0$. In this case, we obtained the following fractional superalgebra :

$$\{Q_a, Q_b, Q_g\} = 0, \tag{31}$$

$$[Q_1, X_1] = Q_2, [Q_2, X_2] = Q_1, [Q_1, X_3] = Q_1, [Q_2, X_3] = -Q_2. \tag{32}$$

Note that, for N=3 the relations (31) are not zero [3].

4. CONCLUSION

By applying fractional super algebra methods which are explained at Ref [3], to $sl(2)$ Lie algebra which is a special case, we obtained N=3 fractional super generalization of $sl(2)$ at n=3. In this generalization some $b_{\alpha\beta\gamma}^j$ structure constants where different from zero.

In this paper, by applying the same method we obtained N=4 fractional super generalization of $sl(2)$ at n=3 and found $b_{\alpha\beta\gamma}^j$ structure constants equal to zero.

REFERENCES

- [1] Raush De Traubenberga, M. and Slupinski, M J 2000 J.Math.Phys 41 4556.
- [2] Raush De Traubenberga, M. , "Fractional supersymmetri and Lie algebras", arXiv:hep-th/0007150.
- [3] Ahmedov, H. , Yıldız, A. and Ucan, Y., "Fractional super Lie algebras and groups", J. Phys. A., 34, 6413, 2001.
- [4] Ahn, C., Bernard, D. and Leclair, A., "Fractional supersymmetries in perturbed coset CFT and integrable solition theory", Nucl. Phys., B 346, 409, 1990.
- [5] Raush De Traubenberga, M. and Slupinski, M.J., Mod. Phys. Lett. A, 39, 3051, 1997.
- [6] Kerner, R., "Z-graded algebras and the cubic root of supersymmetry translations.", J. Math. Phys., 33, 403, 1992 .
- [7] Abramov, V., Kerner, R. and Le Roy, B., "A Z-graded generalization of supersymmetry", J. Math. Phys., 38, 1650, 1997.
- [8] De Azcarraga, J.A. and Macfarlane, M.J., "Group theoretical foundations of fractinal supersymmetry.", J. Math. Phys., 37, 1115, 1996.
- [9] Durand, S., Mod. Phys. Lett., A7, 2905, 1992.
- [10] Ahmedov, H. and Dayı, Ö.F., "Two dimensional fractional supersymmetry from the Quantum Poincare group at roots of unity.", J.Phys. A.32.6247, 1999.
- [11] Vilenkin, N.Ya. and Klimyk, A.U., " Representations of Lie Groups and Special Functions", vol 1, Kluwer Academic Press, The Netherland, 1991.